

# Ground State Solution for a quasilinear coupled system

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In this work we consider the following class of quasilinear coupled systems

$$(S_\theta) \begin{cases} -\Delta u + a(x)u - \Delta(u^2)u = g(u) + \alpha\theta\lambda(x)|u|^{\alpha-2}u|v|^\beta, & x \in \mathbb{R}^N, \\ -\Delta v + b(x)v - \Delta(v^2)v = h(v) + \beta\theta\lambda(x)|v|^{\beta-2}v|u|^\alpha, & x \in \mathbb{R}^N, \end{cases}$$

where  $N \geq 3$  and  $a, b : \mathbb{R}^N \rightarrow \mathbb{R}$  are positive potentials,  $\lambda : \mathbb{R}^N \rightarrow \mathbb{R}$  is a suitable continuous function,  $\theta > 0$  and  $\alpha, \beta > 2$  satisfying  $\alpha + \beta < 2.2^*$ . On the nonlinear terms we assume that  $g, h$  are in  $C^1$  class and are superlinear functions at infinity and at the origin. The main theorem is stated without the well known Ambrosetti-Rabinowitz condition at infinity. Using a change of variable, we turn the quasilinear coupled system into a nonlinear coupled system, where we establish a variational approach based on Nehari method.

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## 1 Introduction

We look for ground states for the general class of quasilinear coupled systems involving Schrödinger equations  $(S_\theta)$ . This class of systems imposes some difficulties. The first one is that the energy functional associated to System  $(S_\theta)$  is not well defined in the whole space  $H^1(\mathbb{R}^N)^2$ . Thus, motivated by seminal works [1, 2, 3, 4, 5, 6] we also use a change of variable to reformulate our initial problem, obtaining a nonlinear coupled system. After change of variable, the modified problem has an associated energy functional well defined in the whole space  $H^1(\mathbb{R}^N)^2$  and the solutions are related with solutions of the initial System  $(S_\theta)$ . The second difficulty is the lack of compactness due to the fact that the system is defined in the whole Euclidean space  $\mathbb{R}^N$ . Moreover, System  $(S_\theta)$  involve strongly coupled Schrödinger equations because of the coupling terms in the right hand side. We emphasize that we do not use the well known Ambrosetti-Rabinowitz condition. We suppose that the potentials  $a, b$  satisfy the following hypotheses:

- (a<sub>0</sub>)  $a, b, \lambda \in C(\mathbb{R}^N, \mathbb{R})$  are 1-periodic functions;
- (a<sub>1</sub>)  $a(x) \geq a_0$  and  $b(x) \geq b_0$  for some  $a_0, b_0 > 0$ ;
- (a<sub>2</sub>)  $\lambda(x) \geq 0$  for all  $x \in \mathbb{R}^N$  and  $\lambda(x) > 0$  for all  $x \in \Omega$ , for some  $\Omega \subset \mathbb{R}^N$  such that  $|\Omega| < +\infty$ ;
- (g<sub>0</sub>)  $g, h \in C^1(\mathbb{R}, \mathbb{R})$ ;
- (g<sub>1</sub>)  $|g(t)| \leq C(1 + |t|^{p-1})$ ,  $|h(t)| \leq C(1 + |t|^{p-1})$ , for all  $t \in \mathbb{R}$  for some  $C > 0$  and  $p \in (4, 2 \cdot 2^*)$ ;
- (g<sub>2</sub>)  $\lim_{t \rightarrow 0} \frac{g(t)}{t} = 0$ ,  $\lim_{t \rightarrow 0} \frac{h(t)}{t} = 0$ ;
- (g<sub>3</sub>)  $\lim_{|t| \rightarrow +\infty} \frac{g(t)}{t^3} = +\infty$ ,  $\lim_{|t| \rightarrow +\infty} \frac{h(t)}{t^3} = +\infty$ ;
- (g<sub>4</sub>) The functions  $t \rightarrow \frac{g(t)}{t^3}$ ,  $t \rightarrow \frac{h(t)}{t^3}$  are strictly increasing in  $|t|$ ;
- (g<sub>5</sub>) There holds  $0 \leq G(t) \leq G(|t|)$  and  $0 \leq H(t) \leq H(|t|)$ , for all  $t \in \mathbb{R}$ .

**Theorem** *Under the above hypothesis, there exists a ground state solution for each  $\theta > 0$ . Moreover there exist  $\theta_0 > 0$  such that the System  $(S_\theta)$  has at least one positive ground state solution, for all  $\theta \geq \theta_0$ .*

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## References

- [1] A. Ambrosetti, C. Colorado, *Bound and ground states of coupled nonlinear Schrödinger equations*. C. R. Math. Acad. Sci. Paris **342** (2006), no. 7, 453-458.
- [2] M. Colin, L. Jeanjean, *Solutions for a quasilinear Schrödinger equation: a dual approach*, Nonlinear Anal. **56** (2004), 213–226.
- [3] J. Liu, Z. Q. Wang, X. Wu, *Multibump solutions for quasilinear elliptic equations with critical growth*. J. Math. Phys. **54** (2013), 121–131.
- [4] J.Q. Liu, Z.Q. Wang, *Soliton solutions for quasilinear Schrödinger equations I*, Proc. Amer. Math. Soc. **131** (2002), 441–448.
- [5] J.Q. Liu, Y.Q. Wang and Z.Q. Wang, *Soliton solutions for quasilinear Schrödinger equations II*, J. Differential Equations **187** (2003), 473–493.
- [6] J.Q. Liu, Y.Q. Wang and Z.Q. Wang, *Solutions for quasilinear Schrödinger equations via the Nehari method*, Comm. Partial Differential Equations **29** (2004), 879–901.
- [7] L. A. Maia, E. Montefusco, B. Pellacci, *Weakly coupled nonlinear Schrödinger systems: the saturation effect*, Calc. Var. Partial Differential Equations **46** (2013), no. 1-2, 325-351.
- [8] L. A. Maia, E. Montefusco, B. Pellacci, *Positive solutions for a weakly coupled nonlinear Schrödinger system*. J. Differential Equations **229** (2006), no. 2, 743-767.
- [9] P.H. Rabinowitz, *On a class of nonlinear Schrödinger equations*, Z. Angew. Math. Phys. **43** (1992), 270–291.
- [10] E.A.B. Silva, G.F. Vieira, *Quasilinear asymptotically periodic Schrödinger equations with subcritical growth*, Nonlinear Analysis **72** (2010) 2935–2949.
- [11] M. Yang, *Existence of solutions for a quasilinear Schrödinger equation with subcritical nonlinearities*, Nonlinear Anal. **75** (2012), 5362–5373.